



# A **Field–Theoretic** Approach to Microtubule Growth

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## Introduction

Reaction-diffusion equations

Master equation

Field-theory

## Field theory of microtubule growth

Model

Tubulin probe in vicinity of MT cap

Discrete microtubule movement

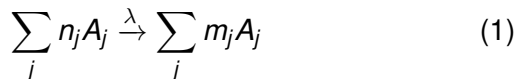
Bath of tubulin

## Conclusion



## Two components

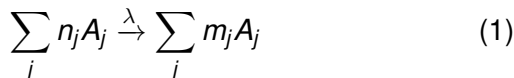
► Chemical reactions





## Two components

- ▶ Chemical reactions



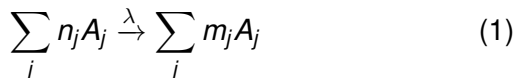
- ▶ Movement: Diffusion, ballistic, anomalous, continuous, discrete

$$\partial_t^\alpha A(t, x) = D(\partial_x^\beta)^2 A(t, x) + V \partial_x A(t, x) \quad (2)$$



## Two components

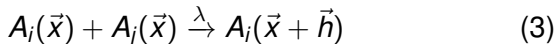
- ▶ Chemical reactions



- ▶ Movement: Diffusion, ballistic, anomalous, continuous, discrete

$$\partial_t^\alpha A(t, \mathbf{x}) = D(\partial_x^\beta)^2 A(t, \mathbf{x}) + V \partial_x A(t, \mathbf{x}) \quad (2)$$

- ▶ Combination of reactions and movement





## ► Lattice space



## Master equation

- ▶ Lattice space
- ▶ discrete particle numbers



## Master equation

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- ▶ reactions are proportional to number of available reactants





- ▶ Lattice space
- ▶ discrete particle numbers
- ▶ reactions are proportional to number of available reactants
- ▶ Example 1: Reaction



$$\begin{aligned} \partial_t \mathcal{P}(\{m\}, \{n\}, t) &= \text{gain} - \text{loss} \\ &= \lambda \sum_x \left( \binom{m_x + 2}{2} \mathcal{P}(\{m_x + 2\}, \{n_x - 1\}, t) + \right. \\ &\quad \left. - \binom{m_x}{2} \mathcal{P}(\{m_x\}, \{n_x\}, t) \right) \end{aligned}$$



► Example 2: Diffusion

$$\partial_t \mathcal{P}(\{m\}, t) = \frac{D}{h^2} \sum_x \sum_{|x-y|=1} \left( (m_x + 1) \mathcal{P}(\{m_x + 1, m_y - 1\}, t) + m_x \mathcal{P}(\{m_x\}, t) \right) \quad (5)$$



## Miraculous transformation (Doi [1], Peliti[2], Wijland[3])

►  $2A \xrightarrow{\lambda} B$

$$H = \int_{\mathbb{R}} \int_{\mathbb{R}^d} \frac{\lambda}{2} \left( \varphi(t, \mathbf{x})^2 \psi^\dagger(t, \mathbf{x}) - \varphi^\dagger(t, \mathbf{x})^2 \varphi(t, \mathbf{x})^2 \right) d^d \mathbf{x} dt \quad (6)$$



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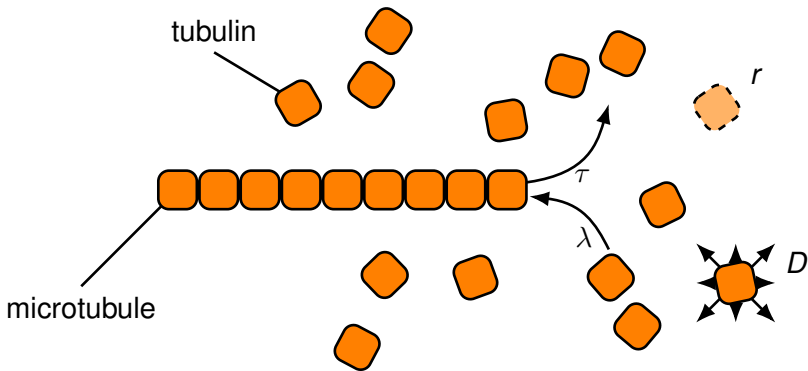
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- ▶ Diffusion

$$H = \int_{\mathbb{R}} \int_{\mathbb{R}^d} D \varphi^\dagger(t, \mathbf{x}) \Delta \varphi(t, \mathbf{x}) d^d \mathbf{x} dt \quad (7)$$



# Simple Model of Microtubule Growth





# How to do field theory in 5 steps

- ▶ Take probabilities of interest:  $\varphi(t, x), \psi(t, x), \dots$



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- ▶ Consider ALL corresponding Feynman diagrams
- ▶ Write their integral representation



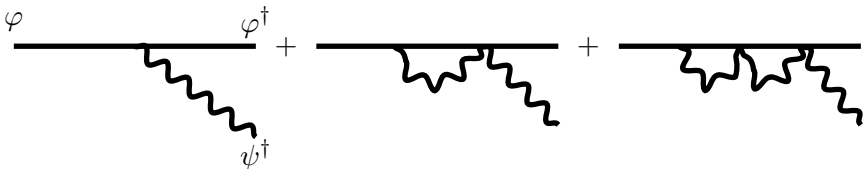
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- ▶ Consider ALL corresponding Feynman diagrams
- ▶ Write their integral representation
- ▶ Calculate...



## How to find single tubulin probability densities

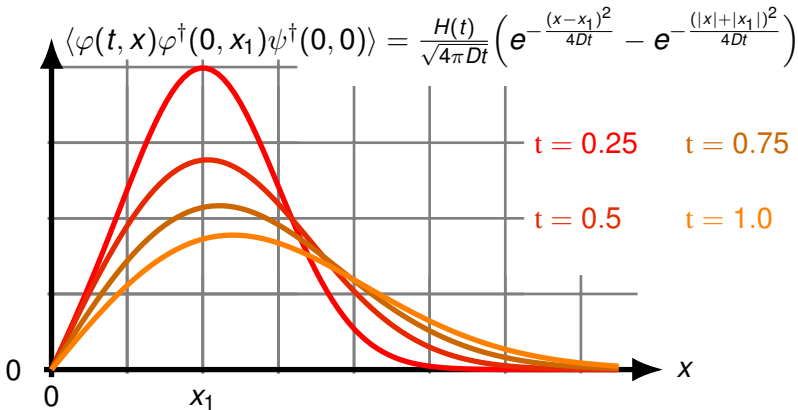
$$\langle \varphi(t, x) \varphi^\dagger(0, x_1) \psi^\dagger(0, 0) \rangle \quad (8)$$





Tubulin probe in vicinity of MT cap

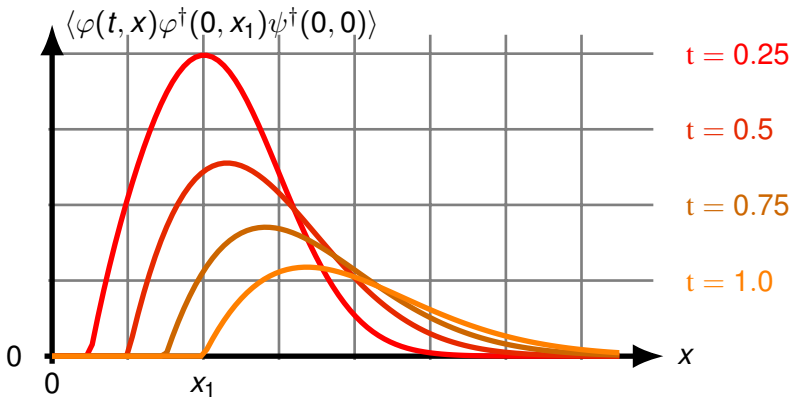
## Single tubulin close to stationary microtubule tip





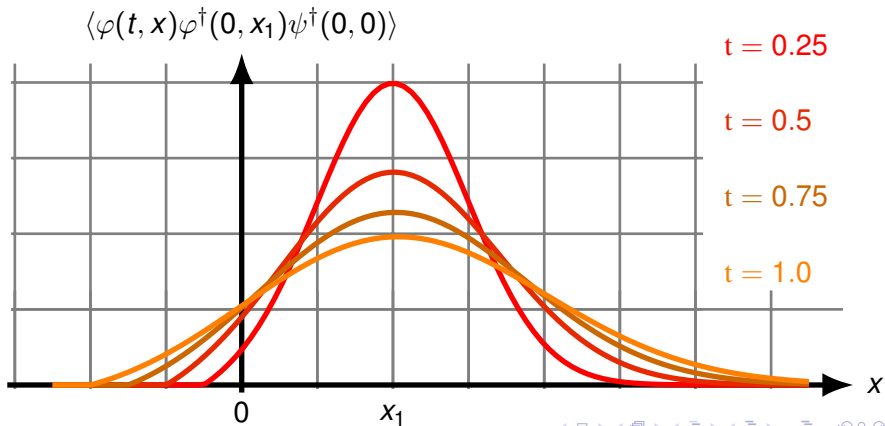
Tubulin probe in vicinity of MT cap

## Single tubulin close to moving microtubule tip I





## Single tubulin close to moving microtubule tip II



## MT cap movement induced by tubulin absorption



$$\langle \psi_j(t) \psi_{j_0}^\dagger(0) \varphi^\dagger(0, x_2) \rangle = \delta_{j, j_0} + \operatorname{erfc}\left(\frac{|x_2 - hj_0|}{\sqrt{4Dt}}\right) (\delta_{j, j_0+1} - \delta_{j, j_0})$$



## How to incorporate large numbers of tubulin?

Homogenous tubulin density  $\zeta$

$$\varphi = \check{\varphi} + \zeta$$

Without decay

$$\langle \psi_j(t) \psi_{j_0}^\dagger(t_0) \rangle_T = \begin{cases} e^{-\lambda\zeta(t-t_0)} \frac{(\lambda\zeta(t-t_0))^{j-j_0}}{(j-j_0)!} & j - j_0 \geq 0 \\ 0 & j - j_0 < 0 \end{cases}$$

With decay

$$\langle \psi_j(t) \psi_{j_0}^\dagger(t_0) \rangle_T = e^{-(\lambda\zeta + \tau)(t-t_0)} \sum_{m=\max\{0, j_0-j\}}^{\infty} \frac{(\tau(t-t_0))^m}{m!} \frac{(\lambda\zeta(t-t_0))^{j-j_0+m}}{(j-j_0+m)!} \quad (9)$$





## Depletion of tubulin

1 dim

$$\langle \varphi(t, \mathbf{x}) \psi_0^\dagger(\mathbf{0}) \rangle_{t \rightarrow \infty} = \zeta + (\tau - \lambda \zeta) \frac{e^{-\sqrt{\frac{r}{D}}|x|}}{2\sqrt{Dr}} \left( 1 - \frac{\lambda}{2\sqrt{Dr} + \lambda} \right)$$

3 dim

$$\langle \varphi(t, \mathbf{x}) \psi_0^\dagger(\mathbf{0}) \rangle_{t \rightarrow \infty} = \zeta + \frac{(\tau - \lambda \zeta) e^{-\sqrt{\frac{r}{D}}|x|}}{4\pi D|x|} \left( 1 - \frac{\lambda(2\Lambda - \pi\sqrt{r/D})}{4\pi^2 D + \lambda(2\Lambda - \pi\sqrt{r/D})} \right)$$



# Take-Home Message

1. Field Theory can produce analytic results for complex systems



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2. From microscopic & short-time behaviour it can predict behaviour in any spatial and time scale






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1. Field Theory can produce analytic results for complex systems
2. From microscopic & short-time behaviour it can predict behaviour in any spatial and time scale
3. It can predict fluctuations and correlations of any order



Thank you!



-  Masao Doi, *Second quantization representation for classical many-particle system*, J. Phys. A: Math. Gen., Vol. 9. No. 9.1976.
-  L. Peliti, *Path integral approach to birth-death processes on a lattice*, Journal de Physique 1985, 46 (9), pp.1469-1483.  
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-  Frédéric van Wijland, *Field theory for reaction-diffusion processes with hard-core particles*, PHYSICAL REVIEW E, VOLUME 63, 022101,2001.